# Search for fractionally charged particles in $p p$ collisions at $\sqrt{s}=7 \mathrm{TeV}$ 

S. Chatrchyan et al. *<br>(CMS Collaboration)

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#### Abstract

A search is presented for free heavy long-lived fractionally charged particles produced in $p p$ collisions at $\sqrt{s}=7 \mathrm{TeV}$. The data sample was recorded by the CMS detector at the LHC and corresponds to an integrated luminosity of $5.0 \mathrm{fb}^{-1}$. Candidate fractionally charged particles are identified by selecting tracks with associated low charge measurements in the silicon tracking detector. Observations are found to be consistent with expectations for background processes. The results of the search are used to set upper limits on the cross section for pair production of fractionally charged, massive spin- $1 / 2$ particles that are neutral under $S U(3)_{C}$ and $S U(2)_{L}$. We exclude at $95 \%$ confidence level such particles with electric charge $\pm 2 e / 3$ with masses below 310 GeV , and those with charge $\pm e / 3$ with masses below 140 GeV .


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## I. INTRODUCTION

The reasons for expecting physics beyond the standard model to manifest itself in $p p$ collisions at the Large Hadron Collider (LHC) are as compelling as ever. As suggested in Ref. [1], one example of physics beyond the standard model that may have eluded previous searches is a new particle with an electric charge less than that of the electron. Owing to their lower ionization energy loss, the trajectories of such fractionally charged particles may not pass typical track quality requirements and a dedicated analysis is required.

While fractionally charged particles are common in some theoretical scenarios (e.g., superstrings [2,3]), a variety of searches for these objects in bulk matter, cosmic rays and accelerator based experiments have reported null results [4]. Strong constraints on models with fractionally charged states come from astrophysics and cosmology [5]. These constraints, however, do not apply if the reheating temperature of the Universe after the last stage of inflation is much lower than the mass of the fractionally charged particle such that there is no thermal relic from freeze-out [6]. We search for such particles, using as a benchmark the scenario considered in [5], namely, a model with new, fractionally charged, massive spin-1/2 particles that are neutral under $S U(3)_{C}$ and $S U(2)_{L}$ and therefore couple only to the photon and the $Z$. We denote such particles with fractional charge $\pm q e$, where $e$ is the charge of the electron, as $L_{q}$, and assume they have a lifetime sufficiently long such that they do not decay within the detector volume.

An interesting possibility is that these $L_{q}$ could also be charged under a new asymptotically free gauge group

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$S U(N)$ with a confinement scale $\Lambda$. This would make them a variant type of "quirk," which are quarklike, naturally fractionally charged particles [7]. For $\Lambda \leqslant 100 \mathrm{eV}$, the confining string between the quirk-antiquirk pair would have a negligible impact on their trajectories over distances typical of collider-detector dimensions. Thus such quirks would have the same kinematic properties and experimental signature as the benchmark model considered in this paper. However, the existence of this string would cause eventual annihilation of any pairs formed in the early Universe resulting in a negligible relic abundance [6], thereby evading the constraints cited by Ref. [5] irrespective of cosmological history.

We search for $L_{q}, \bar{L}_{q}$ particles in a sample of tracks with a muonlike signature. We identify fractionally charged particle candidates by their anomalously low ionization energy loss in the inner tracker. This study complements Compact Muon Solenoid (CMS) searches for heavy stable charged particles with anomalously high ionization energy loss [8,9].

## II. SIGNAL SIMULATION

Pair production of $L_{q} \bar{L}_{q}$ at the LHC proceeds via a modified Drell-Yan mechanism with weak isospin $t_{3 L}=$ 0 , which has $L_{q}-Z$ axial coupling $g_{A}=0$ and vector coupling $g_{V}=-2 q \sin ^{2} \theta_{W}$. We have performed Monte Carlo simulations of this signal using PYTHIA v6.422 [10], with $q=1 / 3,2 / 3$, and 1 , and with masses of $100,200,300$, 400,500 , and 600 GeV . The cross sections are calculated to leading order with the CTEQ6L1 parton distribution functions. The detector response is modeled with a simulation based on GEANT4 [11].

## III. DETECTOR AND DATA SAMPLE

The central feature of the CMS apparatus [12] is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T . Within the field volume are a
silicon pixel and strip tracker, a lead-tungstate crystal electromagnetic calorimeter, and a brass/scintillator hadron calorimeter. Muons are measured in gas-ionization detectors embedded in the steel return yoke. Extensive forward calorimetry complements the coverage provided by the barrel and endcap detectors. Of particular importance to this search is the inner tracker [13], which consists of 1440 silicon pixel and 15148 silicon strip detector modules. The inner tracker records 16 measurements on average per track.

The analysis is performed on the $p p$ collision data sample recorded by the CMS detector at $\sqrt{s}=7 \mathrm{TeV}$ in 2011, corresponding to an integrated luminosity of $5.0 \mathrm{fb}^{-1}$.

The data are selected with a single-muon trigger that requires a track reconstructed in both the inner tracker and muon detectors with transverse momentum $p_{\mathrm{T}}>40 \mathrm{GeV}$ and pseudorapidity $|\eta|<2.1$, where $\eta=-\ln [\tan (\theta / 2)]$ and $\theta$ is the polar angle. The radius of curvature of a fractionally charged particle in a magnetic field is larger than that of a particle of unit charge with the same momentum, so the reconstructed $p_{\mathrm{T}}$ is larger than the true $p_{\mathrm{T}}$ by the inverse of the particle's charge. The trigger requirement that unit charge particles have $p_{\mathrm{T}}>40 \mathrm{GeV}$ corresponds to a requirement of $p_{\mathrm{T}}>27 \mathrm{GeV}$ for $L_{2 / 3}$ and $p_{\mathrm{T}}>13 \mathrm{GeV}$ for $L_{1 / 3}$.

The trigger efficiency for $L_{2 / 3}$ is in the range $67 \%-74 \%$ per event, which is very similar to the efficiency for unit charge particles simulated with the same mass. By contrast, the $L_{1 / 3}$ trigger efficiency is between $8 \%$ and $18 \%$. The lower trigger efficiency for $L_{1 / 3}$ results from the fact that many of the energy deposits in the tracker and muon detectors are below the threshold required to record a measurement. The trigger efficiency for $L_{1 / 3}$ depends on the particle's velocity $\beta$, since $d E / d x \propto 1 / \beta^{2}$ [14]. The reconstruction efficiency for slower-moving $L_{1 / 3}$ particles is larger because the energy deposits are more likely to be above threshold. As a result, the reconstructed velocity distribution is very different from the generated distribution. For $L_{2 / 3}$ the larger overall efficiency means that the velocity distribution is less affected by the reconstruction, but slower-moving $L_{2 / 3}$ particles fail the signal region requirement since their recorded $d E / d x$ measurements are too large.

## IV. SELECTION

In a preselection step, candidates for fractionally charged particles are defined as tracks reconstructed in the inner tracker and matched to a track in the muon detectors. The preselection criteria, which are described below, are chosen to obtain well-reconstructed tracks and to suppress background from cosmic ray muons. We consider candidate muon tracks with large reconstructed transverse momenta, $p_{\mathrm{T}}>45 \mathrm{GeV}$, as measured in the inner
tracker, in the range $|\eta|<1.5$. A loose requirement on the track fit quality, $\chi^{2} /$ dof $<10$, rejects very poorly reconstructed tracks. We also require at least six $d E / d x$ measurements from the tracker, where a $d E / d x$ measurement is the signal amplitude recorded in an inner detector module divided by the particle's path length through the module. To ensure that the track is isolated, the sum of the $p_{\mathrm{T}}$ of all other tracks within a cone of $\Delta R \equiv \sqrt{(\Delta \phi)^{2}+(\Delta \eta)^{2}}<$ 0.3 , where $\phi$ is the azimuthal angle, around the candidate must be less than 0.1 times the $p_{\mathrm{T}}$ of the candidate. The sum of the electromagnetic and hadronic calorimeter energy recorded within this cone, including that deposited by the candidate, must be less than 5 GeV .

Muons from cosmic rays are found in only a small fraction of the triggered events, but because they typically arrive out of coincidence with the bunch crossing, the charge is sampled away from the signal's maximum, and the resulting low $d E / d x$ measurement can be indistinguishable from that of a fractionally charged particle. Several requirements help to suppress the cosmic ray background. The primary vertex with the smallest distance to the point of closest approach of the candidate track is required to be a well-reconstructed vertex [15]. The track must also have at least two $d E / d x$ measurements in the pixel tracker. The candidate track is required to satisfy $\left|d_{z}\right|<0.5 \mathrm{~cm}$ and $\left|d_{x y}\right|<0.1 \mathrm{~cm}$, where $d_{z}$ and $d_{x y}$ are the longitudinal and transverse impact parameters with respect to the primary vertex. Based on the time of flight measurement [9] by the muon detectors, we determine the time the candidate particle was at the interaction point (IP) under the assumptions that the particle had velocity $\beta=1$ and was moving outward from the IP. This time must not be earlier than the nominal bunch crossing that triggered the event, a requirement that rejects over half of the background from $p p$ collisions and cosmic rays. We require $\alpha_{\max }<2.8 \mathrm{rad}$, where $\alpha_{\text {max }}$ is the maximum three-dimensional opening angle between the candidate track and any high-momentum ( $p_{\mathrm{T}}>35 \mathrm{GeV}$ ) track reconstructed in either the inner tracker or in the muon system alone.

The combined efficiency of the trigger and the preselection for signal events generated with mass 100 GeV is $45 \%$ per event for $L_{2 / 3}$ and $4.4 \%$ for $L_{1 / 3}$. The largest efficiency loss is at the trigger stage for $L_{1 / 3}$, where a signal track may fail to be reconstructed. The preselection requirements reduce the signal efficiency by roughly a factor of 2.

After the preselection, events with two $L_{q}$ candidates are rejected from the search sample if the invariant mass of the two candidates $m_{L L}$ is in the range $80<m_{L L}<100 \mathrm{GeV}$. This ensures that the search sample is independent of a control sample, defined below, that is used to estimate the collisions background. Although in the considered signal model, $L_{q} \bar{L}_{q}$ are produced in pairs, events containing a single candidate track are retained after the preselection. Events containing more than two candidates, which make up less than $0.1 \%$ of the search sample, are excluded.

We isolate the signal using a technique that imposes few assumptions on any particular model but nonetheless has the power to suppress the large standard model backgrounds from $p p$ collisions. A fractionally charged particle is most clearly distinguished from a standard model particle by its lower rate of energy loss in the detector since $d E / d x \propto q^{2}$ [14]. Figure 1 shows the distribution of $d E / d x$ measurements associated with tracks passing the preselection, for the search sample, a control sample, and the $100 \mathrm{GeV} L_{2 / 3}$ and $L_{1 / 3}$ simulated signal samples. The measured values from data lie predominantly in the region above $2 \mathrm{MeV} / \mathrm{cm}$. We therefore define low-ionizing measurements to be those with $d E / d x<2 \mathrm{MeV} / \mathrm{cm}$. By requiring a number of such low $d E / d x$ measurements, standard model backgrounds can be suppressed while most of the signal events that pass the preselection are retained. Tracks that intersect a sensor close to its edge or near the boundary between two sensor modules can result in low $d E / d x$ measurements because of the partial collection of the deposited charge, so these track measurements are not considered in the analysis. The distance to the sensor edge for which $d E / d x$ measurements are excluded is between $0.5 \%$ and $5 \%$ of the distance to the center of the module, depending on the tracker subsystem.

A signal region is determined by maximizing the expected mass limit on the production cross section of fractionally charged particles while varying the minimum


FIG. 1 (color online). Distribution of $d E / d x$ measurements associated with tracks passing the preselection in the search sample and the $Z$-peak control sample. Simulated $L_{q} \bar{L}_{q}$ signal samples for a mass of 100 GeV are shown for $q=2 / 3,1 / 3$. The distributions are normalized to the area of the search sample. The magenta vertical line at $d E / d x=2 \mathrm{MeV} / \mathrm{cm}$ indicates the upper limit of the range of measurements considered to be low ionizing. The inset is an enlargement of the region of low $d E / d x$, plotted on a semilogarithmic scale.
number of low $d E / d x$ measurements. The signal region optimization is performed for simulated samples with a $L_{q}$ mass of 100 GeV and 400 GeV , and for both samples the optimum signal region is defined with the requirement that events contain a track with at least six low $d E / d x$ measurements. For the 100 GeV signal events passing the preselection, $75 \%$ of $L_{2 / 3}$ events are in the signal region, and $93 \%$ of $L_{1 / 3}$ events are in the signal region.

## V. BACKGROUND ESTIMATE

We use control samples of data to estimate the background contribution from cosmic ray muons and from particles produced in $p p$ collisions. The data-driven method provides a background estimate without the use of simulation.

To estimate the cosmic ray background, we use a sample of muons obtained with the nominal preselection except for two inverted requirements, $0.1<\left|d_{x y}\right|<1.1 \mathrm{~cm}$ and $0.5<\left|d_{z}\right|<50 \mathrm{~cm}$. The yield in the signal region of this sample is scaled by a weight factor to obtain the cosmic ray background for the nominal preselection. This weight factor is the product of two weights, each determined from a cosmic ray enriched sample as the ratio of the yield in the nominal selection region to the yield after inverting a single requirement. This scaled yield gives a cosmic ray background estimate of 0.007 events.

A Z-peak control sample is used to estimate the $p p$ collision background. This sample is selected by relaxing the transverse momentum requirement to $p_{\mathrm{T}}>35 \mathrm{GeV}$, requiring $80<m_{L L}<100 \mathrm{GeV}$, and applying all other preselection requirements. Figure 1 shows the distributions of $d E / d x$ measurements associated with selected tracks for both the search sample and the Z-peak control sample. The two distributions agree within the statistical uncertainties over the full $d E / d x$ range. The ratio of the number of tracks passing the preselection in the search sample to the number of tracks passing the preselection in the Z-peak control sample is 10.5 . The control sample is scaled by this ratio to model the distribution in the search sample.

The simulation of the control sample is also shown in Fig. 1. This simulation is used only to assess the uncertainty in the signal efficiency, since the background estimate is entirely data driven. The inset in Fig. 1 is an enlargement of the region of low $d E / d x$, plotted on a semilogarithmic scale. This inset also shows the results of a modified simulation, which includes the effect of a possible source of anomalously low $d E / d x$ hits not reproduced in the nominal simulation. The background simulations are discussed in the next section.

To estimate the background in the signal region, we extrapolate from the background-dominated region of events containing a preselected track with zero to five low $d E / d x$ measurements. For a muon from a $Z$ decay, the $d E / d x$ measurements associated with the track are
expected to be uncorrelated, and the number of measurements below a given $d E / d x$ value can be described by a generalized binomial function,

$$
\begin{aligned}
& N_{\mathrm{evts}}=N_{0}\binom{\mu}{n} p^{n}(1-p)^{\mu-n} \\
& \binom{\mu}{n}=\frac{\Gamma(\mu+1)}{\Gamma(n+1) \Gamma(\mu-n+1)}
\end{aligned}
$$

where $N_{\text {evts }}$ is the number of events containing at least one track with $n$ low $d E / d x$ measurements, $\mu$ is the average number of measurements per track, $p$ is the probability for a single measurement to be low $d E / d x, N_{0}$ is a normalization factor, and $\Gamma(n)$ is the gamma function. The fit of the binomial function to the background-dominated region of the Z-peak control sample is shown in Fig. 2. The fitted parameters are $\mu=17.5 \pm 1.7, p=0.0125 \pm 0.0013$, and $N_{0}=(5.03 \pm 0.03) \times 10^{6}$; the values of $\mu$ and $p$ are close to those expected based on the number of measurements per track and the fraction of low $d E / d x$ measurements. This function describes the distribution in the control sample well, with $\chi^{2} /$ dof $=0.07 / 1$, corresponding to a $\chi^{2}$ probability of $79 \%$. This is strong support for the hypothesis that the data are distributed binomially and therefore that the $d E / d x$ measurements are uncorrelated.


FIG. 2 (color online). Number of events with at least one track with the given number of low $d E / d x$ measurements, for search data and the scaled $Z$-peak control sample background estimate. The binomial function fit to the control sample is shown, with the band representing its uncertainty. The ratio of the data to the binomial fit is also presented. No tracks in the search sample have five or more low $d E / d x$ measurements.

Extrapolation of the fitted binomial function into the signal region yields a $p p$ background estimate of 0.005 events.

## VI. SYSTEMATIC UNCERTAINTIES

The systematic uncertainties that significantly impact the results are the uncertainties in the integrated luminosity, the background estimate, and the signal efficiency. The uncertainty in the integrated luminosity is $2.2 \%$ [16].

The cosmic ray background estimate has a statistical uncertainty of $71 \%$ that arises from the relatively small size of the sample with inverted $d_{x y}$ and $d_{z}$ requirements used for its determination. The statistical uncertainties in the weighting factors are $1 \%$ and $24 \%$ for the $d_{x y}$ and $d_{z}$ requirements, respectively. The systematic uncertainty associated with the assumption that the $d_{x y}$ and $d_{z}$ variables are uncorrelated is assessed by examining a sample defined by replacing the inverted $d_{z}$ selection with an inverted $\alpha_{\text {max }}$ requirement. This sample, obtained by requiring $0.1<\left|d_{x y}\right|<1.1 \mathrm{~cm}, \alpha_{\text {max }}>2.8 \mathrm{rad}$, and all other preselection criteria, provides a second estimate of the cosmic ray background, which differs from the nominal estimate by $42 \%$. The statistical and systematic uncertainties are summed in quadrature; the total cosmic ray background estimate is $0.007 \pm 0.006$ events.

We assess three potential sources of uncertainty in the $p p$ background estimate in the signal region. The first source is from the choice of the function used to fit the control sample. While this is often a large source of uncertainty in many a posteriori fits to data, our hypothesis that a binomial function describes the distribution of the number of low $d E / d x$ measurements is motivated a priori from first principles. We do not expect a large uncertainty from this source. For completeness, however, other functions are also compared to the data. Fits of several modified exponential, power-law, and polynomial functions fail to converge or have very low $\chi^{2}$ probabilities. One function that does fit the distribution reasonably well is $N_{\text {evts }}=$ $p_{0} n^{p_{1}+p_{2} n}$, where $p_{i}$ are free parameters. The difference between the background estimate from this function and the nominal background estimate is 0.001 events.

The second potential source of uncertainty in the $p p$ background estimate arises from the statistical uncertainties in the fitted parameters of the binomial function. The propagation of these uncertainties results in an uncertainty in the background estimate of $\pm 0.0004$ events.

A third source of uncertainty arises from the small disagreement between the distribution of low $d E / d x$ measurements from the control sample and that from the search sample. In the background-dominated region, the largest statistically significant discrepancy between the two samples is $9 \%$, for zero low $d E / d x$ measurements. To assess the resulting systematic uncertainty, the control sample fit is repeated for a large number of trials, in each case setting the value of the distribution in each bin randomly, according to a Gaussian distribution with a mean of the original value
and width of $9 \%$ of the original value. The RMS of the background estimates from all of these trials is 0.004 events, which is taken as the uncertainty due to the accuracy with which the control sample models the search sample.

The three sources of uncertainty in the $p p$ background estimate are summed in quadrature giving a total estimate of $0.005 \pm 0.004$ events. The high precision of this estimate is due to the large statistics of the control sample, which leads to small statistical uncertainties in the fitted parameters. Likewise, the high degree of accuracy of the background estimate is reflected in the relatively small systematic uncertainty assigned. This is a direct consequence of the somewhat unusual aspect of this analysis that the functional form with which the data would be distributed under the background-only hypothesis was derived from first principles.

The sum of the $p p$ and cosmic ray backgrounds gives a total background estimate of $0.012 \pm 0.007$ events.

This search uses the data themselves to estimate all backgrounds, so the uncertainties in the simulation only impact the determination of the efficiency of the benchmark signal model. The systematic uncertainties in the signal efficiencies are summarized in Table I.

The simulation of the trigger efficiency for a fractionally charged particle is sensitive to the accurate modeling of the muon detectors' electronics and gas gain as well as the threshold for recording a hit because it has less ionization energy loss, and thus the peak of its Landau distribution is closer to the threshold than that of a muon. We assess the impact on the signal region efficiency of a variation in the simulated gain of the muon system by a conservative estimate of its uncertainty. The impact of such a variation on $L_{2 / 3}$ is small, since the charge distributions are typically far above the threshold. However, for $L_{1 / 3}$, the probability to record a hit degrades significantly as the gain decreases. The impact of this variation for $L_{2 / 3}$ is $1 \%$, and for $L_{1 / 3}$ is $8.5 \%$. The systematic uncertainty in the offline global muon identification requirement is negligible by comparison.

The modeling of the tracker $d E / d x$ measurements impacts the simulated signal efficiency by affecting the efficiency of track reconstruction and signal region selection. Larger tracker energy deposits are more likely to be above the threshold required to record a measurement, and the track reconstruction efficiency increases with more measurements. Larger $d E / d x$ measurements also reduce the

TABLE I. Systematic uncertainties (in \%) in the signal efficiency.

| Source | $L_{1 / 3}$ | $L_{2 / 3}$ |
| :--- | :---: | :---: |
| Muon trigger | 8.5 | 1 |
| Tracker $d E / d x$ measurements | 16 | 7.5 |
| Track momentum scale | $<1$ | $<1$ |
| Muon timing | 2 | 2 |
| Total | 18 | 8 |

fraction of reconstructed tracks that are in the signal region, since fewer measurements are below the $2 \mathrm{MeV} / \mathrm{cm}$ limit. To evaluate the accuracy of the simulation of the $d E / d x$ measurements, we compare the $d E / d x$ distributions of the Z-peak control sample in simulation and in data, as shown in Fig. 1. The agreement in the low $d E / d x$ region is evaluated as the shift of all $d E / d x$ measurements required to obtain the same fraction below $2 \mathrm{MeV} / \mathrm{cm}$ in both samples. For the nominal selection, the simulated $d E / d x$ distribution must be shifted by $2.6 \%$ to obtain the same fraction below $2 \mathrm{MeV} / \mathrm{cm}$ as in the data. For a larger sample obtained with looser selection criteria, the required shift is $5 \%$. We use the larger of these values, $5 \%$, as an estimate of the agreement between simulation and data. To assess the resulting uncertainty in the signal efficiency, we vary the amplitude of the $d E / d x$ measurements by $\pm 5 \%$ before resimulating the trigger emulation, track reconstruction, and full selection. A variation of $+5 \%$ in the $d E / d x$ measurements changes the signal efficiency by $+16 \%$ for $L_{1 / 3}$ particles and $-7.5 \%$ for $L_{2 / 3}$. The efficiency changes in opposite directions because for $L_{1 / 3}$ the increased reconstruction efficiency is the dominant effect, while for $L_{2 / 3}$ the smaller signal region efficiency has a greater impact. These variations in the signal efficiency are taken as the systematic uncertainties associated with the modeling of the tracker $d E / d x$ measurements.

Potential causes of incorrect modeling of $d E / d x$ in our simulation that could produce a disagreement at low $d E / d x$ have been examined. The most likely possibility is a residual source of low $d E / d x$ hits that are not removed by the sensor-edge fiducial cuts. Such a source could be

TABLE II. Numbers of events in the signal region, containing a track with at least six $d E / d x$ measurements with $d E / d x<$ $2 \mathrm{MeV} / \mathrm{cm}$, estimated from the background and observed in the search data.

| Cosmic rays | $0.007 \pm 0.006$ |
| :--- | :---: |
| $p p$ collisions | $0.005 \pm 0.004$ |
| Total background | $0.012 \pm 0.007$ |
| Observed events | 0 |

TABLE III. The signal efficiency and average velocity $\langle\beta\rangle$ of events in the signal region for different mass points and charge hypotheses.

| Mass <br> $(\mathrm{GeV})$ | $L_{2 / 3}$ <br> Signal eff. | $\langle\beta\rangle$ | Signal eff. | $\langle\beta\rangle$ |
| :--- | :---: | :---: | :---: | :---: |
| 100 | $0.341 \pm 0.026$ | 0.84 | $0.041 \pm 0.007$ | 0.52 |
| 200 | $0.357 \pm 0.027$ | 0.83 | $0.060 \pm 0.011$ | 0.51 |
| 300 | $0.337 \pm 0.026$ | 0.82 | $0.074 \pm 0.013$ | 0.51 |
| 400 | $0.314 \pm 0.024$ | 0.80 | $0.091 \pm 0.016$ | 0.51 |
| 500 | $0.265 \pm 0.020$ | 0.79 | $0.104 \pm 0.019$ | 0.51 |
| 600 | $0.251 \pm 0.019$ | 0.78 | $0.109 \pm 0.019$ | 0.51 |



FIG. 3 (color online). Expected and observed limits on the cross section of $L_{q} \bar{L}_{q}$ for $q=2 / 3$ and $1 / 3$. The theoretical prediction of modified Drell-Yan production with $t_{3 L}=0$ is shown. The lines for the expected and observed limits are overlapping.
accommodated by the control sample data if, at most, $0.06 \%$ of all measurements are affected. A simulation assuming a mismeasurement rate at this level reproduces the observed data distribution, as shown in the inset of Fig. 1. Such an effect would impact less than $1 \%$ of all tracks and change the signal efficiency by an even smaller amount. Furthermore, the likelihood of a track to have six such anomalous measurements is extremely small, so the effect on the background estimate would be negligible.

The uncertainty associated with the track momentum scale is less than $1 \%$. The impact of the uncertainty in the muon timing measurements is $2 \%$. This is assessed by varying the timing measurements according to the measured discrepancy between the data and simulation, as described in [9].

## VII. RESULTS

The numbers of expected background and observed events are summarized in Table II. We observe zero events in the data search sample in the signal region, which is consistent with the background estimate of $0.012 \pm 0.007$ events. The $L_{q} \bar{L}_{q}$ signal efficiency and average $\beta$ for different signal masses and charges are listed in Table III. Ninety-five percent confidence level (CL) upper limits on the $L_{q} \bar{L}_{q}$ production cross section are calculated using the $\mathrm{CL}_{\mathrm{s}}$ criterion [17,18]. Expected and observed 95\% CL limits are shown in Fig. 3. These limits vary from 1.7 to
2.3 fb , for $q=2 / 3$, and from 14 to 5.4 fb , for $q=1 / 3$, for masses between 100 and 600 GeV . We exclude the production of $L_{2 / 3}$ with a mass below 310 GeV and the production of $L_{1 / 3}$ with a mass below 140 GeV at $95 \% \mathrm{CL}$.

## VIII. CONCLUSION

A search has been performed for heavy, long-lived, leptonlike fractionally charged particles, using the signature of at least six low $d E / d x$ measurements. The search is based on a $p p$ collision sample recorded by the CMS detector at $\sqrt{s}=7 \mathrm{GeV}$, corresponding to an integrated luminosity of $5.0 \mathrm{fb}^{-1}$. Zero events are observed in the signal region, consistent with the background estimate. Upper limits are set with $95 \%$ confidence on the cross section of pair produced, spin-1/2 particles that are neutral under $S U(3)_{C}$ and $S U(2)_{L}$. The existence of such particles with masses less than 310 GeV is excluded by these cross section limits for the case $q=2 / 3$, and with masses less than 140 GeV for $q=1 / 3$.

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S. Costa, ${ }^{55 a, 55 b}$ R. Potenza, ${ }^{55 a, 55 b}$ A. Tricomi,,${ }^{55 a, 55 b}$ C. Tuve,,${ }^{55 a, 55 b}$ G. Barbagli, ${ }^{56 a}$ V. Ciulli, ${ }^{56 a, 56 b}$ C. Civinini, ${ }^{56 a}$ R. D'Alessandro, ${ }^{56 a, 56 b}$ E. Focardi, ${ }^{56 a, 56 b}$ S. Frosali, ${ }^{56 a, 56 b}$ E. Gallo, ${ }^{56 \mathrm{a}}$ S. Gonzi, ${ }^{56,56 b}$ M. Meschini, ${ }^{56 \mathrm{a}}$ S. Paoletti, ${ }^{56 \mathrm{a}}$ G. Sguazzoni, ${ }^{56 \mathrm{a}}$ A. Tropiano, ${ }^{56 \mathrm{a}} \mathrm{L}$. Benussi, ${ }^{57}$ S. Bianco, ${ }^{57}$ S. Colafranceschi, ${ }^{57, \mathrm{z}}$ F. Fabbri, ${ }^{57}$ D. Piccolo, ${ }^{57}$
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Stephans, ${ }^{139}$ F. Stöckli, ${ }^{139}$ K. Sumorok, ${ }^{139}$ K. Sung, ${ }^{139}$ D. Velicanu, ${ }^{139}$ E. A. Wenger, ${ }^{139}$ R. Wolf, ${ }^{139}$ B. Wyslouch, ${ }^{139}$ M. Yang, ${ }^{139}$ Y. Yilmaz, ${ }^{139}$ A. S. Yoon, ${ }^{139}$ M. Zanetti, ${ }^{139}$ S. I. Cooper, ${ }^{140}$ B. Dahmes, ${ }^{140}$ A. De Benedetti, ${ }^{140}$ G. Franzoni, ${ }^{140}$ A. Gude, ${ }^{140}$ S. C. Kao, ${ }^{140}$ K. Klapoetke, ${ }^{140}$ Y. Kubota, ${ }^{140}$ J. Mans, ${ }^{140}$ N. Pastika, ${ }^{140}$ R. Rusack, ${ }^{140}$ M. Sasseville, ${ }^{140}$ A. Singovsky, ${ }^{140}$ N. Tambe, ${ }^{140}$ J. Turkewitz, ${ }^{140}$ L. M. Cremaldi, ${ }^{141}$ R. Kroeger, ${ }^{141}$ L. Perera, ${ }^{141}$ R. Rahmat, ${ }^{141}$ D. A. Sanders, ${ }^{141}$ E. Avdeeva, ${ }^{142}$ K. Bloom,,$^{142}$ S. Bose, ${ }^{142}$ J. Butt, ${ }^{142}$ D. R. Claes, ${ }^{142}$ A. Dominguez, ${ }^{142}$ M. Eads, ${ }^{142}$ J. Keller, ${ }^{142}$ I. Kravchenko, ${ }^{142}$ J. Lazo-Flores, ${ }^{142}$ H. Malbouisson, ${ }^{142}$ S. Malik, ${ }^{142}$ G. R. Snow, ${ }^{142}$ A. Godshalk, ${ }^{143}$ I. Iashvili, ${ }^{143}$ S. Jain, ${ }^{143}$ A. Kharchilava, ${ }^{143}$ A. Kumar, ${ }^{143}$ G. Alverson, ${ }^{144}$ E. Barberis, ${ }^{144}$ D. Baumgartel, ${ }^{144}$ M. Chasco, ${ }^{144}$ J. Haley, ${ }^{144}$ D. Nash, ${ }^{144}$ D. Trocino, ${ }^{144}$ D. Wood, ${ }^{144}$ J. Zhang, ${ }^{144}$ A. Anastassov, ${ }^{145}$ A. Kubik, ${ }^{145}$ N. Mucia, ${ }^{145}$ N. Odell, ${ }^{145}$ R. A. Ofierzynski, ${ }^{145}$ B. Pollack, ${ }^{145}$ A. Pozdnyakov, ${ }^{145}$ M. Schmitt, ${ }^{145}$ S. Stoynev, ${ }^{145}$ M. Velasco, ${ }^{145}$ S. Won, ${ }^{145}$ L. Antonelli, ${ }^{146}$ D. Berry, ${ }^{146}$ A. Brinkerhoff, ${ }^{146}$ K. M. Chan, ${ }^{146}$ M. Hildreth, ${ }^{146}$ C. Jessop, ${ }^{146}$ D. J. Karmgard, ${ }^{146}$ J. Kolb, ${ }^{146}$ K. Lannon, ${ }^{146}$ W. Luo, ${ }^{146}$ S. Lynch, ${ }^{146}$ N. Marinelli, ${ }^{146}$ D. M. Morse, ${ }^{146}$ T. Pearson, ${ }^{146}$ M. Planer, ${ }^{146}$ R. Ruchti, ${ }^{146}$ J. Slaunwhite, ${ }^{146}$ N. Valls, ${ }^{146}$ M. Wayne, ${ }^{146}$ M. Wolf, ${ }^{146}$ B. Bylsma, ${ }^{147}$ L. S. Durkin, ${ }^{147}$ C. Hill, ${ }^{147}$ R. Hughes, ${ }^{147}$ K. Kotov, ${ }^{147}$ T. Y. Ling, ${ }^{147}$ D. Puigh, ${ }^{147}$ M. Rodenburg, ${ }^{147}$ C. Vuosalo, ${ }^{147}$ G. Williams, ${ }^{147}$ B. L. Winer, ${ }^{147}$ N. Adam, ${ }^{148}$ E. Berry, ${ }^{148}$ P. Elmer, ${ }^{148}$ D. Gerbaudo, ${ }^{148}$ V. Halyo, ${ }^{148}$ P. Hebda, ${ }^{148}$ J. Hegeman, ${ }^{148}$ A. Hunt, ${ }^{148}$ P. Jindal, ${ }^{148}$ D. Lopes Pegna, ${ }^{148}$ P. Lujan, ${ }^{148}$ D. Marlow, ${ }^{148}$ T. Medvedeva, ${ }^{148}$ M. Mooney, ${ }^{148}$ J. Olsen, ${ }^{148}$ P. Piroué, ${ }^{148}$ X. Quan, ${ }^{148}$ A. Raval,,${ }^{148}$ B. Safdi, ${ }^{148}$ H. Saka, ${ }^{148}$ D. Stickland, ${ }^{148}$ C. Tully, ${ }^{148}$ J. S. Werner, ${ }^{148}$ A. Zuranski, ${ }^{148}$ E. Brownson, ${ }^{149}$ A. Lopez, ${ }^{149}$ H. Mendez, ${ }^{149}$ J. E. Ramirez Vargas, ${ }^{149}$ E. Alagoz, ${ }^{150}$ V. E. Barnes, ${ }^{150}$ D. Benedetti, ${ }^{150}$ G. Bolla, ${ }^{150}$ D. Bortoletto, ${ }^{150}$ M. De Mattia, ${ }^{150}$ A. Everett, ${ }^{150}$ Z. Hu, ${ }^{150}$ M. Jones, ${ }^{150}$ O. Koybasi, ${ }^{150}$ M. Kress, ${ }^{150}$ A. T. Laasanen, ${ }^{150}$ N. Leonardo, ${ }^{150}$ V. Maroussov, ${ }^{150}$ P. Merkel, ${ }^{150}$ D. H. Miller, ${ }^{150}$ N. Neumeister, ${ }^{150}$ I. Shipsey, ${ }^{150}$ D. Silvers, ${ }^{150}$ A. Svyatkovskiy, ${ }^{150}$ M. Vidal Marono, ${ }^{150}$ H. D. Yoo,,${ }^{150}$ J. Zablocki, ${ }^{150}$ Y. Zheng, ${ }^{150}$ S. Guragain,,$^{151}$ N. Parashar, ${ }^{151}$ A. Adair, ${ }^{152}$ C. Boulahouache, ${ }^{152}$ K. M. Ecklund, ${ }^{152}$ F. J. M. Geurts, ${ }^{152}$ W. Li, ${ }^{152}$ B. P. Padley, ${ }^{152}$ R. Redjimi, ${ }^{152}$ J. Roberts, ${ }^{152}$ J. Zabel, ${ }^{152}$ B. Betchart, ${ }^{153}$ A. Bodek, ${ }^{153}$ Y. S. Chung, ${ }^{153}$ R. Covarelli, ${ }^{153}$ P. de Barbaro, ${ }^{153}$ R. Demina, ${ }^{153}$ Y. Eshaq, ${ }^{153}$ T. Ferbel, ${ }^{153}$ A. Garcia-Bellido, ${ }^{153}$ P. Goldenzweig, ${ }^{153}$ J. Han, ${ }^{153}$ A. Harel, ${ }^{153}$ D. C. Miner, ${ }^{153}$ D. Vishnevskiy, ${ }^{153}$ M. Zielinski, ${ }^{153}$ A. Bhatti, ${ }^{154}$ R. Ciesielski, ${ }^{154}$ L. Demortier, ${ }^{154}$ K. Goulianos, ${ }^{154}$ G. Lungu, ${ }^{154}$ S. Malik, ${ }^{154}$ C. Mesropian, ${ }^{154}$ S. Arora, ${ }^{155}$ A. Barker, ${ }^{155}$ J. P. Chou, ${ }^{155}$ C. Contreras-Campana, ${ }^{155}$ E. Contreras-Campana, ${ }^{155}$ D. Duggan, ${ }^{155}$ D. Ferencek, ${ }^{155}$ Y. Gershtein, ${ }^{155}$ R. Gray, ${ }^{155}$ E. Halkiadakis, ${ }^{155}$ D. Hidas, ${ }^{155}$ A. Lath, ${ }^{155}$ S. Panwalkar, ${ }^{155}$ M. Park, ${ }^{155}$ R. Patel, ${ }^{155}$ V. Rekovic, ${ }^{155}$ J. Robles, ${ }^{155}$ K. Rose, ${ }^{155}$ S. Salur, ${ }^{155}$ S. Schnetzer, ${ }^{155}$ C. Seitz, ${ }^{155}$ S. Somalwar, ${ }^{155}$ R. Stone, ${ }^{155}$ S. Thomas, ${ }^{155}$ G. Cerizza, ${ }^{156}$ M. Hollingsworth, ${ }^{156}$ S. Spanier, ${ }^{156}$ Z. C. Yang, ${ }^{156}$ A. York, ${ }^{156}$ R. Eusebi, ${ }^{157}$ W. Flanagan, ${ }^{157}$ J. Gilmore, ${ }^{157}$ T. Kamon, ${ }^{157, g g g}$ V. Khotilovich, ${ }^{157}$ R. Montalvo, ${ }^{157}$ I. Osipenkov, ${ }^{157}$ Y. Pakhotin, ${ }^{157}$ A. Perloff, ${ }^{157}$ J. Roe, ${ }^{157}$
A. Safonov, ${ }^{157}$ T. Sakuma, ${ }^{157}$ S. Sengupta, ${ }^{157}$ I. Suarez, ${ }^{157}$ A. Tatarinov, ${ }^{157}$ D. Toback, ${ }^{157}$ N. Akchurin, ${ }^{158}$ J. Damgov, ${ }^{158}$ C. Dragoiu, ${ }^{158}$ P. R. Dudero, ${ }^{158}$ C. Jeong, ${ }^{158}$ K. Kovitanggoon, ${ }^{158}$ S. W. Lee, ${ }^{158}$ T. Libeiro, ${ }^{158}$ Y. Roh, ${ }^{158}$ I. Volobouev, ${ }^{158}$ E. Appelt, ${ }^{159}$ A. G. Delannoy, ${ }^{159}$ C. Florez, ${ }^{159}$ S. Greene, ${ }^{159}$ A. Gurrola, ${ }^{159}$ W. Johns, ${ }^{159}$ P. Kurt, ${ }^{159}$ C. Maguire, ${ }^{159}$ A. Melo, ${ }^{159}$ M. Sharma, ${ }^{159}$ P. Sheldon, ${ }^{159}$ B. Snook, ${ }^{159}$ S. Tuo, ${ }^{159}$ J. Velkovska, ${ }^{159}$
M. W. Arenton, ${ }^{160}$ M. Balazs, ${ }^{160}$ S. Boutle, ${ }^{160}$ B. Cox, ${ }^{160}$ B. Francis, ${ }^{160}$ J. Goodell, ${ }^{160}$ R. Hirosky, ${ }^{160}$ A. Ledovskoy, ${ }^{160}$ C. Lin, ${ }^{160}$ C. Neu, ${ }^{160}$ J. Wood, ${ }^{160}$ S. Gollapinni, ${ }^{161}$ R. Harr, ${ }^{161}$ P. E. Karchin, ${ }^{161}$ C. Kottachchi Kankanamge Don, ${ }^{161}$ P. Lamichhane, ${ }^{161}$ A. Sakharov, ${ }^{161}$ M. Anderson, ${ }^{162}$ D. Belknap, ${ }^{162}$ L. Borrello, ${ }^{162}$ D. Carlsmith, ${ }^{162}$ M. Cepeda, ${ }^{162}$ S. Dasu, ${ }^{162}$ E. Friis, ${ }^{162}$ L. Gray, ${ }^{162}$ K. S. Grogg, ${ }^{162}$ M. Grothe, ${ }^{162}$ R. Hall-Wilton, ${ }^{162}$ M. Herndon, ${ }^{162}$ A. Hervé, ${ }^{162}$ P. Klabbers, ${ }^{162}$ J. Klukas, ${ }^{162}$ A. Lanaro, ${ }^{162}$ C. Lazaridis, ${ }^{162}$ J. Leonard, ${ }^{162}$ R. Loveless, ${ }^{162}$ A. Mohapatra, ${ }^{162}$ I. Ojalvo, ${ }^{162}$ F. Palmonari, ${ }^{162}$ G. A. Pierro, ${ }^{162}$ I. Ross, ${ }^{162}$ A. Savin, ${ }^{162}$ W. H. Smith, ${ }^{162}$ and J. Swanson ${ }^{162}$
(CMS Collaboration)

${ }^{1}$ Yerevan Physics Institute, Yerevan, Armenia<br>${ }^{2}$ Institut für Hochenergiephysik der OeAW, Wien, Austria<br>${ }^{3}$ National Centre for Particle and High Energy Physics, Minsk, Belarus<br>${ }^{4}$ Universiteit Antwerpen, Antwerpen, Belgium<br>${ }^{5}$ Vrije Universiteit Brussel, Brussel, Belgium<br>${ }^{6}$ Université Libre de Bruxelles, Bruxelles, Belgium<br>${ }^{7}$ Ghent University, Ghent, Belgium<br>${ }^{8}$ Université Catholique de Louvain, Louvain-la-Neuve, Belgium<br>${ }^{9}$ Université de Mons, Mons, Belgium<br>${ }^{10}$ Centro Brasileiro de Pesquisas Fisicas, Rio de Janeiro, Brazil<br>${ }^{11}$ Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil<br>${ }^{12}$ Instituto de Fisica Teorica, Universidade Estadual Paulista, Sao Paulo, Brazil<br>${ }^{13}$ Institute for Nuclear Research and Nuclear Energy, Sofia, Bulgaria<br>${ }^{14}$ University of Sofia, Sofia, Bulgaria<br>${ }^{15}$ Institute of High Energy Physics, Beijing, China<br>${ }^{16}$ State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China<br>${ }^{17}$ Universidad de Los Andes, Bogota, Colombia<br>${ }^{18}$ Technical University of Split, Split, Croatia<br>${ }^{19}$ University of Split, Split, Croatia<br>${ }^{20}$ Institute Rudjer Boskovic, Zagreb, Croatia<br>${ }^{21}$ University of Cyprus, Nicosia, Cyprus<br>${ }^{22}$ Charles University, Prague, Czech Republic<br>${ }^{23}$ Academy of Scientific Research and Technology of the Arab Republic of Egypt,<br>Egyptian Network of High Energy Physics, Cairo, Egypt<br>${ }^{24}$ National Institute of Chemical Physics and Biophysics, Tallinn, Estonia<br>${ }^{25}$ Department of Physics, University of Helsinki, Helsinki, Finland<br>${ }^{26}$ Helsinki Institute of Physics, Helsinki, Finland<br>${ }^{27}$ Lappeenranta University of Technology, Lappeenranta, Finland<br>${ }^{28}$ DSM/IRFU, CEA/Saclay, Gif-sur-Yvette, France<br>${ }^{29}$ Laboratoire Leprince-Ringuet, Ecole Polytechnique, IN2P3-CNRS, Palaiseau, France<br>${ }^{30}$ Institut Pluridisciplinaire Hubert Curien, Université de Strasbourg, Université de Haute Alsace Mulhouse, CNRS/IN2P3, Strasbourg, France<br>${ }^{31}$ Centre de Calcul de l'Institut National de Physique Nucleaire et de Physique des Particules, CNRS/IN2P3, Villeurbanne, France, Villeurbanne, France<br>${ }^{32}$ Université de Lyon, Université Claude Bernard Lyon 1, CNRS-IN2P3, Institut de Physique Nucléaire de Lyon, Villeurbanne, France<br>${ }^{33}$ Institute of High Energy Physics and Informatization, Tbilisi State University, Tbilisi, Georgia<br>${ }^{34}$ RWTH Aachen University, I. Physikalisches Institut, Aachen, Germany<br>${ }^{35}$ RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany<br>${ }^{36}$ RWTH Aachen University, III. Physikalisches Institut B, Aachen, Germany<br>${ }^{37}$ Deutsches Elektronen-Synchrotron, Hamburg, Germany<br>${ }^{38}$ University of Hamburg, Hamburg, Germany<br>${ }^{39}$ Institut für Experimentelle Kernphysik, Karlsruhe, Germany<br>${ }^{40}$ Institute of Nuclear Physics "Demokritos", Aghia Paraskevi, Greece<br>${ }^{41}$ University of Athens, Athens, Greece

${ }^{42}$ University of Ioánnina, Ioánnina, Greece<br>${ }^{43}$ KFKI Research Institute for Particle and Nuclear Physics, Budapest, Hungary<br>${ }^{44}$ Institute of Nuclear Research ATOMKI, Debrecen, Hungary<br>${ }^{45}$ University of Debrecen, Debrecen, Hungary<br>${ }^{46}$ Panjab University, Chandigarh, India<br>${ }^{47}$ University of Delhi, Delhi, India<br>${ }^{48}$ Saha Institute of Nuclear Physics, Kolkata, India<br>${ }^{49}$ Bhabha Atomic Research Centre, Mumbai, India<br>${ }^{50}$ Tata Institute of Fundamental Research-EHEP, Mumbai, India<br>${ }^{51}$ Tata Institute of Fundamental Research-HECR, Mumbai, India<br>${ }^{52}$ Institute for Research in Fundamental Sciences (IPM), Tehran, Iran<br>${ }^{53 a}$ INFN Sezione di Bari, Bari, Italy<br>${ }^{53 \mathrm{~b}}$ Università di Bari, Bari, Italy<br>${ }^{53 \mathrm{c}}$ Politecnico di Bari, Bari, Italy<br>${ }^{54 \mathrm{a}}$ INFN Sezione di Bologna, Bologna, Italy<br>${ }^{54 \mathrm{~b}}$ Università di Bologna, Bologna, Italy<br>${ }^{55 a}$ INFN Sezione di Catania, Catania, Italy<br>${ }^{55 \mathrm{~b}}$ Università di Catania, Catania, Italy<br>${ }^{56 \mathrm{a}}$ INFN Sezione di Firenze, Firenze, Italy<br>${ }^{56 \mathrm{~b}}$ Università di Firenze, Firenze, Italy<br>${ }^{57}$ INFN Laboratori Nazionali di Frascati, Frascati, Italy<br>${ }^{58 a}$ INFN Sezione di Genova, Genova, Italy<br>${ }^{58 \mathrm{~b}}$ Università di Genova, Genova, Italy<br>${ }^{59 \mathrm{a}}$ INFN Sezione di Milano-Bicocca, Milano, Italy<br>${ }^{59 \mathrm{~b}}$ Università di Milano-Bicocca, Milano, Italy<br>${ }^{60 \mathrm{a}}$ INFN Sezione di Napoli, Napoli, Italy<br>${ }^{60 \mathrm{~b}}$ Università di Napoli "Federico II", Napoli, Italy<br>${ }^{61 \mathrm{a}}$ INFN Sezione di Padova, Padova, Italy<br>${ }^{61 \mathrm{~b}}$ Università di Padova, Padova, Italy<br>${ }^{61 \mathrm{c}}$ Università di Trento (Trento), Padova, Italy<br>${ }^{62 \mathrm{a}}$ INFN Sezione di Pavia, Pavia, Italy<br>${ }^{62 \mathrm{~b}}$ Università di Pavia, Pavia, Italy<br>${ }^{63 \mathrm{a}}$ INFN Sezione di Perugia, Perugia, Italy<br>${ }^{63 \mathrm{~b}}$ Università di Perugia, Perugia, Italy<br>${ }^{64 \mathrm{a}}$ INFN Sezione di Pisa, Pisa, Italy<br>${ }^{64 \mathrm{~b}}$ Università di Pisa, Pisa, Italy<br>${ }^{64 \mathrm{c}}$ Scuola Normale Superiore di Pisa, Pisa, Italy<br>${ }^{65 \mathrm{a}}$ INFN Sezione di Roma, Roma, Italy ${ }^{65 \mathrm{~b}}$ Università di Roma, Roma, Italy<br>${ }^{66 a}$ INFN Sezione di Torino, Torino, Italy<br>${ }^{66 \mathrm{~b}}$ Università di Torino, Torino, Italy<br>${ }^{66 \mathrm{c}}$ Università del Piemonte Orientale (Novara), Torino, Italy<br>${ }^{67}$ INFN Sezione di Trieste, Trieste, Italy<br>${ }^{67 \mathrm{~b}}$ Università di Trieste, Trieste, Italy<br>${ }^{68}$ Kangwon National University, Chunchon, Korea<br>${ }^{69}$ Kyungpook National University, Daegu, Korea<br>${ }^{70}$ Chonnam National University, Institute for Universe and Elementary Particles, Kwangju, Korea<br>${ }^{71}$ Korea University, Seoul, Korea<br>${ }^{72}$ University of Seoul, Seoul, Korea<br>${ }^{73}$ Sungkyunkwan University, Suwon, Korea<br>${ }^{74}$ Vilnius University, Vilnius, Lithuania<br>${ }^{75}$ Centro de Investigacion y de Estudios Avanzados del IPN, Mexico City, Mexico<br>${ }^{76}$ Universidad Iberoamericana, Mexico City, Mexico<br>${ }^{77}$ Benemerita Universidad Autonoma de Puebla, Puebla, Mexico<br>${ }^{78}$ Universidad Autónoma de San Luis Potosí, San Luis Potosí, Mexico<br>${ }^{79}$ University of Auckland, Auckland, New Zealand<br>${ }^{80}$ University of Canterbury, Christchurch, New Zealand<br>${ }^{81}$ National Centre for Physics, Quaid-I-Azam University, Islamabad, Pakistan<br>${ }^{82}$ National Centre for Nuclear Research, Swierk, Poland<br>${ }^{83}$ Institute of Experimental Physics, Faculty of Physics, University of Warsaw, Warsaw, Poland<br>${ }^{84}$ Laboratório de Instrumentação e Física Experimental de Partículas, Lisboa, Portugal

${ }^{85}$ Joint Institute for Nuclear Research, Dubna, Russia<br>${ }^{86}$ Petersburg Nuclear Physics Institute, Gatchina (St. Petersburg), Russia<br>${ }^{87}$ Institute for Nuclear Research, Moscow, Russia<br>${ }^{88}$ Institute for Theoretical and Experimental Physics, Moscow, Russia<br>${ }^{89}$ Moscow State University, Moscow, Russia<br>${ }^{90}$ P. N. Lebedev Physical Institute, Moscow, Russia<br>${ }^{91}$ State Research Center of Russian Federation, Institute for High Energy Physics, Protvino, Russia<br>${ }^{92}$ University of Belgrade, Faculty of Physics and Vinca Institute of Nuclear Sciences, Belgrade, Serbia<br>${ }^{93}$ Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain<br>${ }^{94}$ Universidad Autónoma de Madrid, Madrid, Spain<br>${ }^{95}$ Universidad de Oviedo, Oviedo, Spain<br>${ }^{96}$ Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, Santander, Spain<br>${ }^{97}$ CERN, European Organization for Nuclear Research, Geneva, Switzerland<br>${ }^{98}$ Paul Scherrer Institut, Villigen, Switzerland<br>${ }^{99}$ Institute for Particle Physics, ETH Zurich, Zurich, Switzerland<br>${ }^{100}$ Universität Zürich, Zurich, Switzerland<br>${ }^{101}$ National Central University, Chung-Li, Taiwan<br>${ }^{102}$ National Taiwan University (NTU), Taipei, Taiwan<br>${ }^{103}$ Chulalongkorn University, Bangkok, Thailand<br>${ }^{104}$ Cukurova University, Adana, Turkey<br>${ }^{105}$ Middle East Technical University, Physics Department, Ankara, Turkey<br>${ }^{106}$ Bogazici University, Istanbul, Turkey<br>${ }^{107}$ Istanbul Technical University, Istanbul, Turkey<br>${ }^{108}$ National Scientific Center, Kharkov Institute of Physics and Technology, Kharkov, Ukraine<br>${ }^{109}$ University of Bristol, Bristol, United Kingdom<br>${ }^{110}$ Rutherford Appleton Laboratory, Didcot, United Kingdom<br>${ }^{111}$ Imperial College, London, United Kingdom<br>${ }^{112}$ Brunel University, Uxbridge, United Kingdom<br>${ }^{113}$ Baylor University, Waco, Texas, USA<br>${ }^{114}$ The University of Alabama, Tuscaloosa, Alabama, USA<br>${ }^{115}$ Boston University, Boston, Massachusetts, USA<br>${ }^{116}$ Brown University, Providence, Rhode, Island, USA<br>${ }^{117}$ University of California, Davis, Davis, California, USA<br>${ }^{118}$ University of California, Los Angeles, Los Angeles, California, USA<br>${ }^{119}$ University of California, Riverside, Riverside, California, USA<br>${ }^{120}$ University of California, San Diego, La Jolla, USA<br>${ }^{121}$ University of California, Santa Barbara, Santa Barbara, California, USA<br>${ }^{122}$ California Institute of Technology, Pasadena, California, USA<br>${ }^{123}$ Carnegie Mellon University, Pittsburgh, Pennsylvania, USA<br>${ }^{124}$ University of Colorado at Boulder, Boulder, Colorado, USA<br>${ }^{125}$ Cornell University, Ithaca, New York, USA<br>${ }^{126}$ Fairfield University, Fairfield, Connecticut, USA<br>${ }^{127}$ Fermi National Accelerator Laboratory, Batavia, Illinois, USA<br>${ }^{128}$ University of Florida, Gainesville, Florida, USA<br>${ }^{129}$ Florida International University, Miami, Florida, USA<br>${ }^{130}$ Florida State University, Tallahassee, Florida, USA<br>${ }^{131}$ Florida Institute of Technology, Melbourne, Florida, USA<br>${ }^{132}$ University of Illinois at Chicago (UIC), Chicago, Illinois, USA<br>${ }^{133}$ The University of Iowa, Iowa City, Iowa, USA<br>${ }^{134}$ Johns Hopkins University, Baltimore, Maryland, USA<br>${ }^{135}$ The University of Kansas, Lawrence, Kansas, USA<br>${ }^{136}$ Kansas State University, Manhattan, Kansas, USA<br>${ }^{137}$ Lawrence Livermore National Laboratory, Livermore, California, USA<br>${ }^{138}$ University of Maryland, College Park, Maryland, USA<br>${ }^{139}$ Massachusetts Institute of Technology, Cambridge, Massachusetts, USA<br>${ }^{140}$ University of Minnesota, Minneapolis, USA<br>${ }^{141}$ University of Mississippi, Oxford, Mississippi, USA<br>${ }^{142}$ University of Nebraska-Lincoln, Lincoln, Nebraska, USA<br>${ }^{143}$ State University of New York at Buffalo, Buffalo, New York, USA<br>${ }^{144}$ Northeastern University, Boston, Massachusetts, USA<br>${ }^{145}$ Northwestern University, Evanston, Illinois, USA

${ }^{146}$ University of Notre Dame, Notre Dame, Indiana, USA<br>${ }^{147}$ The Ohio State University, Columbus, Ohio, USA<br>${ }^{148}$ Princeton University, Princeton, New Jersey, USA<br>${ }^{149}$ University of Puerto Rico, Mayaguez, Puerto, Rico, USA<br>${ }^{150}$ Purdue University, West Lafayette, Indiana, USA<br>${ }^{151}$ Purdue University Calumet, Hammond, Indiana, USA<br>${ }^{152}$ Rice University, Houston, Texas, USA<br>${ }^{153}$ University of Rochester, Rochester, New York, USA<br>${ }^{154}$ The Rockefeller University, New York, New York, USA<br>${ }^{155}$ Rutgers, the State University of New Jersey, Piscataway, New Jersey, USA<br>${ }^{156}$ University of Tennessee, Knoxville, Tennessee, USA<br>${ }^{157}$ Texas A\&M University, College Station, Texas, USA<br>${ }^{158}$ Texas Tech University, Lubbock, Texas, USA<br>${ }^{159}$ Vanderbilt University, Nashville, Tennessee, USA<br>${ }^{160}$ University of Virginia, Charlottesville, Virginia, USA<br>${ }^{161}$ Wayne State University, Detroit, Michigan, USA<br>${ }^{162}$ University of Wisconsin, Madison, Wisconsin, USA

${ }^{a}$ Deceased.
${ }^{\mathrm{b}}$ Also at Vienna University of Technology, Vienna, Austria.
${ }^{\text {c }}$ Also at National Institute of Chemical Physics and Biophysics, Tallinn, Estonia.
${ }^{\mathrm{d}}$ Also at Universidade Federal do ABC, Santo Andre, Brazil.
${ }^{\mathrm{e}}$ Also at California Institute of Technology, Pasadena, CA, USA.
${ }^{\mathrm{f}}$ Also at CERN, European Organization for Nuclear Research, Geneva, Switzerland.
${ }^{\mathrm{g}}$ Also at Laboratoire Leprince-Ringuet, Ecole Polytechnique, IN2P3-CNRS, Palaiseau, France.
${ }^{\text {h }}$ Also at Suez Canal University, Suez, Egypt.
${ }^{\text {i }}$ Also at Zewail City of Science and Technology, Zewail, Egypt.
${ }^{\mathrm{j}}$ Also at Cairo University, Cairo, Egypt.
${ }^{\mathrm{k}}$ Also at Fayoum University, El-Fayoum, Egypt.
${ }^{1}$ Also at British University, Cairo, Egypt.
${ }^{m}$ Now at Ain Shams University, Cairo, Egypt.
${ }^{\mathrm{n}}$ Also at National Centre for Nuclear Research, Swierk, Poland.
${ }^{\circ}$ Also at Université de Haute-Alsace, Mulhouse, France.
${ }^{\mathrm{p}}$ Now at Joint Institute for Nuclear Research, Dubna, Russia.
${ }^{\mathrm{q}}$ Also at Moscow State University, Moscow, Russia.
${ }^{\mathrm{r}}$ Also at Brandenburg University of Technology, Cottbus, Germany.
${ }^{\text {s}}$ Also at Institute of Nuclear Research ATOMKI, Debrecen, Hungary.
${ }^{\text {t}}$ Also at Eötvös Loránd University, Budapest, Hungary.
${ }^{\text {u }}$ Also at Tata Institute of Fundamental Research-HECR, Mumbai, India.
${ }^{\mathrm{v}}$ Also at University of Visva-Bharati, Santiniketan, India.
${ }^{w}$ Also at Sharif University of Technology, Tehran, Iran.
${ }^{x}$ Also at Isfahan University of Technology, Isfahan, Iran.
${ }^{y}$ Also at Plasma Physics Research Center, Science and Research Branch, Islamic Azad University, Tehran, Iran.
${ }^{\text {z }}$ Also at Facoltà Ingegneria Università di Roma, Roma, Italy.
${ }^{\text {aa }}$ Also at Università della Basilicata, Potenza, Italy.
${ }^{\text {bb }}$ Also at Università degli Studi Guglielmo Marconi, Roma, Italy.
${ }^{\mathrm{cc}}$ Also at Università degli Studi di Siena, Siena, Italy.
${ }^{\mathrm{dd}}$ Also at University of Bucharest, Faculty of Physics, Bucuresti-Magurele, Romania.
${ }^{\mathrm{ee}}$ Also at Faculty of Physics of University of Belgrade, Belgrade, Serbia.
${ }^{\text {ff }}$ Also at University of California, Los Angeles, Los Angeles, CA, USA.
${ }^{\text {gq }}$ Also at Scuola Normale e Sezione dell' INFN, Pisa, Italy.
${ }^{\text {hh }}$ Also at INFN Sezione di Roma, Università di Roma, Roma, Italy.
${ }^{\text {ii }}$ Also at University of Athens, Athens, Greece.
${ }^{\mathrm{jj}}$ Also at Rutherford Appleton Laboratory, Didcot, United Kingdom.
${ }^{\text {kk }}$ Also at The University of Kansas, Lawrence, KS, USA.
${ }^{11}$ Also at Paul Scherrer Institut, Villigen, Switzerland.
${ }^{\mathrm{mm}}$ Also at Institute for Theoretical and Experimental Physics, Moscow, Russia.
${ }^{\mathrm{nn}}$ Also at Gaziosmanpasa University, Tokat, Turkey.
${ }^{00}$ Also at Adiyaman University, Adiyaman, Turkey.
${ }^{\mathrm{pp}}$ Also at Izmir Institute of Technology, Izmir, Turkey.
${ }^{\mathrm{qq}}$ Also at The University of Iowa, Iowa City, IA, USA.
${ }^{\text {rr }}$ Also at Mersin University, Mersin, Turkey.
${ }^{\text {ss }}$ Also at Ozyegin University, Istanbul, Turkey.
"Also at Kafkas University, Kars, Turkey.
${ }^{\text {uu }}$ Also at Suleyman Demirel University, Isparta, Turkey.
${ }^{{ }^{v v}}$ Also at Ege University, Izmir, Turkey.
${ }^{w w}$ Also at School of Physics and Astronomy, University of Southampton, Southampton, United Kingdom.
${ }^{x x}$ Also at INFN Sezione di Perugia, Università di Perugia, Perugia, Italy.
${ }^{\text {yy }}$ Also at University of Sydney, Sydney, Australia.
${ }^{\mathrm{zz}}$ Also at Utah Valley University, Orem, UT, USA.
${ }^{\text {aaa }}$ Also at Institute for Nuclear Research, Moscow, Russia.
${ }^{\text {bbb }}$ Also at University of Belgrade, Faculty of Physics and Vinca Institute of Nuclear Sciences, Belgrade, Serbia.
${ }^{c c c}$ Also at Argonne National Laboratory, Argonne, IL, USA.
${ }^{\text {ddd }}$ Also at Erzincan University, Erzincan, Turkey.
${ }^{\text {eee }}$ Also at Mimar Sinan University, Istanbul, Istanbul, Turkey.
${ }^{\text {fff }}$ Also at KFKI Research Institute for Particle and Nuclear Physics, Budapest, Hungary.
${ }^{g g g}$ Also at Kyungpook National University, Daegu, Korea.


[^0]:    *Full author list given at the end of the article.

